BOSON-CONTROLLED QUANTUM TRANSPORT



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Topic: Quantum particle strongly interacting with a correlated/fluctuating background medium



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in many cases motion bears resemblance to the "Echternacher Springprozession" \circleon

related publication: A. Alvermann, D. M. Edwards, HF, Phys. Rev. Lett. 88, 056602 ('07)

(D) (A) (B) (B)



MOTIVATION I

- Strongly correlated charge transport in doped Mott insulators:
 - high-T_c cuprates (AFM spin background)



classical spins: "string effect" hole is bound to its starting point *quantum spins:* "fluctuations" spin lattice can heal itself with rate controlled by exchange parameter

 \rightsquigarrow t – J– type models

colossal magnetoresistive manganites (FM spin background)



strong Hund's rule coupling \sim double-exchange model

in addition: orbital anisotropy of hopping & EP (JT) coupling

Spin/orbital degrees of freedom might be represented by (e.g. Schwinger) bosons!



MOTIVATION II

- Charge transport in systems coupled to phonon or bath degrees of freedom
 - ▶ polarons/excitons, also in CDW materials, DNA,... (deformable lattice)



polaron motion - diagonal *vs* nondiagonal transitions (band *vs* hopping transport)

 \sim Holstein-, Fröhlich-, SSH- or Peierls-Hubbard-type models

▶ low-D systems – nanowires, quantum dots (disorder, phonons, T > 0)



system, contacts/leads, bath,...

appropriate "microscopic" description/modelling?

Again the transport is strongly boson affected, maybe even fluctuation-induced, but now the correlations within the "background" might be week or even absentle



How to capture this great variety of transport phenomena in a simplified model?

Let's consider the following rather general (spinless!) Hamiltonian:

$$H = -t_b \sum_{\langle i,j \rangle} c_j^{\dagger} c_i (b_i^{\dagger} + b_j) - \lambda \sum_i (b_i^{\dagger} + b_i) + \omega_0 \sum_i b_i^{\dagger} b_i + \frac{N\lambda^2}{\omega_0}$$

hopping boson relaxation boson energy

Electron emits or absorbs a local boson every time it hops between lattice sites [but hopping creates (destroys) a boson only on the site the particle leaves (enters)!]:

$$\begin{split} & \mathsf{R}_{i} = \mathsf{c}_{i+1}^{\dagger} \mathsf{c}_{i} \mathsf{b}_{i}^{\dagger} \quad \left| \cdot \quad \stackrel{\frown}{\odot} \quad \cdot \right\rangle \quad \mapsto \quad \left| \cdot \quad \star \quad \odot \right\rangle \\ & \mathsf{L}_{i} = \mathsf{c}_{i-1}^{\dagger} \mathsf{c}_{i} \mathsf{b}_{i}^{\dagger} \quad \left| \cdot \quad \stackrel{\frown}{\odot} \quad \cdot \right\rangle \quad \mapsto \quad \left| \odot \quad \star \quad \cdot \right\rangle \\ & \mathsf{L}_{i}^{\dagger} = \mathsf{c}_{i}^{\dagger} \mathsf{c}_{i-1} \mathsf{b}_{i} \quad \left| \stackrel{\frown}{\odot} \quad \star \quad \cdot \right\rangle \quad \mapsto \quad \left| \cdot \quad \odot \quad \cdot \right\rangle \\ & \mathsf{R}_{i}^{\dagger} = \mathsf{c}_{i}^{\dagger} \mathsf{c}_{i+1} \mathsf{b}_{i} \quad \left| \cdot \quad \star \quad \stackrel{\frown}{\odot} \right\rangle \quad \mapsto \quad \left| \cdot \quad \odot \quad \cdot \right\rangle \end{split}$$

• $\lambda = 0$ - model is analogous to the classical spin model \rightsquigarrow "string effect"? • $\lambda > 0$ allows a boson to decay spontaneously \rightsquigarrow healing of the "spin lattice"



• Note that " $\mathbf{R}_{i}^{(6)} = \mathbf{L}_{i+2}^{\dagger} \mathbf{L}_{i+1}^{\dagger} \mathbf{R}_{i}^{\dagger} \mathbf{L}_{i+2} \mathbf{R}_{i+1} \mathbf{R}_{i}^{*}$ " acts as " $\mathbf{c}_{i+2}^{\dagger} \mathbf{c}_{i}^{*}$: $|\widehat{\circ} \cdot \cdot\rangle \rightarrow |\star \ \widehat{\circ} \cdot \rangle \rightarrow |\star \ \widehat{\circ} \rangle \rightarrow |\star \ \widehat{\circ} \rangle \rightarrow |\star \ \widehat{\circ} \rangle \rightarrow |\widehat{\circ} \ \star \rangle \rightarrow |\widehat{\circ} \ \star \rangle \rightarrow |\widehat{\circ} \ \star \rangle \rightarrow |\cdot \ \widehat{\circ} \ \star \rangle \rightarrow |\cdot \ \widehat{\circ} \rightarrow |\cdot \rangle \rightarrow |\cdot \ \widehat{\circ} \rightarrow |\cdot \ \widehat{\circ} \rightarrow |\cdot \rangle \rightarrow |\cdot \ \widehat{\circ} \rightarrow |\cdot \rangle \rightarrow |\cdot \circ \rightarrow |\cdot \rangle \rightarrow |\cdot \circ \rightarrow |\cdot \rangle \rightarrow |\cdot \rangle$

 \sim lowest order vacuum-restoring process: 1D analogue of 2D "Trugman path" !

 \bullet Unitary transformation $b_i \mapsto b_i + t_f/2t_b$ of H

$$H' = -t_f \sum_{\langle i,j\rangle} c_j^{\dagger} c_i \ -t_b \sum_{\langle i,j\rangle} c_j^{\dagger} c_i (b_i^{\dagger} + b_j) \ + \omega_0 \sum_i b_i^{\dagger} b_i$$

• Different from the t-J model physics of $H^{(\,\prime\,)}$ is governed by *two* energy ratios: t_b/t_f and t_b/ω_0 , where $t_f=2\lambda t_b/\omega_0!$

Obviously H' (H) captures the interplay of "coherent" and "incoherent" transport channels realized in many condensed matter systems!



What does it mean: "Solution"?

• Ground state properties

Adapt variational Hilbert space construction developed for the Holstein/JT polaron problem (see, e.g., Ku, Trugman, Bonča: PRB 65, 174306 ('02)):



 $|1\rangle~e^-$ at site 0 with no phonon excitation

 $|2\rangle~e^-$ and phonon at site 0

 $|3\rangle$ e^- at site 1 and one phonon at site 0

i.e., vertical bonds create or destroy phonons

act m times with off-diagonal terms + all translations on an infinite lattice

One- (two-) particle sector: In most cases 10^4-10^6 basis states are sufficient to obtain an 8-16 digit accuracy for E_0 , $\langle 0| \dots |0\rangle$, ... in any dimension! Note that E_0 calculated this way is variational for the *infinite* system!

 Spectral properties at T=0, thermodynamics Employ Kernel Polynomial Method designed for high-resolution applications: resolution ∝ 1/number of Chebyshev moments! (history goes back 40 years, for a recent review see Rev. Mod. Phys. 78, 275 ('06))



- ground state energy E_0 , kinetic energy part $E_{kin} = \langle 0|H \omega_0 \sum_i b_i^{\dagger} b_i |0\rangle$
- quasiparticle band dispersion E(k), effective mass $1/m^* = \frac{\partial^2 E(k)}{\partial k^2}|_{k=0}$
- particle-boson correlation function $\chi_{ij}=\langle 0|b_i^{\dagger}b_ic_j^{\dagger}c_j|0\rangle$
- one-particle spectral function $A(k,\omega) = \sum_n |\langle n | c_k^{\dagger} | vac \rangle|^2 \, \delta[\omega \omega_n]$
- optical conductivity $\operatorname{Re}\sigma(\omega) = 2\pi D \,\delta(\omega) + \sigma_{\operatorname{reg}}(\omega)$,

regular part $\sigma_{\mathsf{reg}}(\omega) = \pi \sum_{n>0} \frac{|\langle n|j|0\rangle|^2}{\omega_n} \left[\delta(\omega - \omega_n) + \delta(\omega + \omega_n) \right]$,

 $\begin{array}{ll} \text{where } j=j_f+j_b \quad \text{with} \quad j_f=it_f\sum_i \, c_{i+1}^\dagger c_i - c_i^\dagger c_{i+1} \\ j_b=it_b\sum_i \, c_{i+1}^\dagger c_i b_i^\dagger - c_i^\dagger c_{i+1} b_i - c_{i-1}^\dagger c_i b_i^\dagger + c_i^\dagger c_{i-1} b_i \\ \end{array}$

• f-sum rule: $\int_{-\infty}^{\infty} \sigma(\omega) d\omega = 2\pi D + 2 \int_{0}^{\infty} \sigma_{\text{reg}}(\omega) d\omega = -\pi E_{\text{kin}}$

 \sim consistency check: Drude weight $D = 1/2m^*$ (Kohn's formula) \checkmark



• D scaled to the kinetic energy:



 \blacktriangleright free particle: $t_b=0 \rightsquigarrow D=t_f,$ i.e., $-D/E_{kin}=0.5$

▶ weight of lowest order (vacuum restoring) process scales as t_b^6/ω_0^5 \sim boson assisted transport dominates for large $(t_b/\omega_0)^5(t_b/t_f)$

$$\blacktriangleright$$
 D at $\mathrm{t_f}=0$ saturates for $\omega_0
ightarrow 0$





$t_f \leqslant t_b$ (ω_0 not too small):

- pronounced NN particle-boson correlations
- strongly renormalised but well-defined quasiparticle band (reminiscent of spin polaron in the t-J model)
- \bullet optical response threshold given by ω_0
- $\sigma^{reg} \simeq \sigma_b^{reg}$ $S_{tot}(\omega) = \int_0^{\omega} \sigma^{reg}(\omega') d\omega'$
- A(k, ω) signals coherent transport
- → "collective" particle-boson dynamics!

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Limit $t_f = 0 \ (\lambda = 0)$



- exact numerical solution
- particle is still itinerant (but D is small)
- incoherent contributions
- $k \rightarrow k + \pi$ symmetry

► m-boson analytical solution Green function decomposition technique ~> (matrix) continued fraction representation

 $m \leqslant 3$ exact solution possible; only a finite number of states is accessible for the infinite system

 $m \ge 4$ infinitely many states will survive





$t_f \gg t_b$ (ω_0 rather small):

- bosons form a cloud around the particle but are not further correlated
- band flattening near the Brillouin zone boundary

Both is reminiscent of large lattice polarons e.g. in the Holstein model!

- optical response broad absorption feature
- overdamped character of $A(k\omega)$ near $k = 0, \pi$
- system is almost transparent at $k=\pi/2$

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 \sim "diffusive" transport!





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SUMMARY

We studied the interplay of collective dynamics and damping in the presence of correlations and fluctuations within a newly proposed transport model.

The model covers basic aspects of very different Hamiltonians: Hubbard, $t - J \dots$, Fröhlich, Holstein, \dots , SSH - type.

Exact numerical solution $(N \rightarrow \infty) \sim$ surprisingly rich physics:

- moving particle creates local distortions of substantial energy in the medium, which may be able to relax
- their relaxation rate determines how fast the particle can move
- "free" particle \Leftarrow magnetic polaron \Leftarrow lattice polaron
- coherent (correlated) = incoherent (diffusive) transport
- bosonic fluctuations act in two competing ways: limit transport & assist transport!

And all this is obtained for just one particle! Plus background!



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OPEN PROBLEMS I

• Quantum phase transition at half-filling?



- "metal-insulator" transition as λ (t_f) decreases...
- band structure reflects strong correlations...



- Influence of spatial dimensionality?
- What about two particles binding?
- Finite-density effects (0 $\leqslant n \leqslant 1)?$
- Finite-temperature effects?
- Different lattice structures? Frustration!

VED ✓ VED ✓ DMRG (?) QMC ?

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 We need a better analytical understanding of the model, maybe at least for some important limiting cases!
 [1D, λ = 0 & 3-4 boson approximation,... more formal derivation of the model, semiclassical limit (?),...]

Everybody is invited to contribute...